

# **Constructive Models of Discrete and Continuous Physical Phenomena**

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## Discrete Physical Phenomena

- Collisions of rigid objects
- Switching in electrical circuits
- Actuation from software controllers

Ultimately, all of these are processes in the continuum of the physical world. But that is rarely the best way to model them.

## The Model vs. the Thing Being Modeled



Solomon Wolf Golomb (1932) mathematician and engineer and a professor of electrical engineering at the University of Southern California.

## The Model vs. the Thing Being Modeled

*You will never strike oil by  
drilling through the map!*



Solomon Wolf Golomb (1932) mathematician and engineer and a professor of electrical engineering at the University of Southern California.

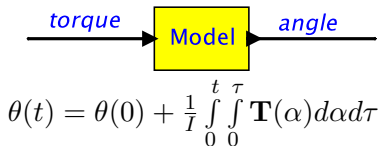
**The value of models depends on the fidelity to what is modeled.**

Physical System:



Image: Wikimedia Commons

Model:



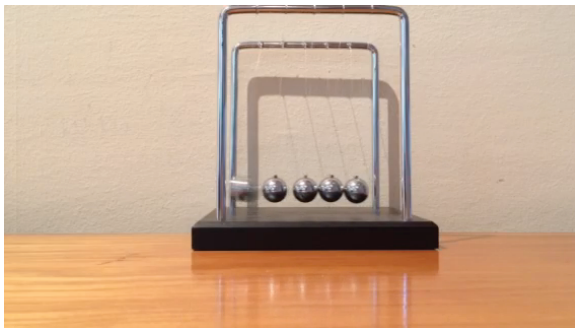
Differential equations are commonly used to model physical dynamics. Such models are always approximate, but often useful. But they have difficulty with discrete events, such as collisions.

## Newton's Cradle – The Model



Image by Dominique Toussaint, [GNU Free Documentation License](#), Version 1.2 or later.

## Newton's Cradle – A Physical Realization



## Model Fidelity

- In *science*, the thing being modeled is given.  
It is up to the model to be faithful.
- In *engineering*, the thing being modeled is a human construction.  
It bears some responsibility to be faithful to the model.

In *engineering*, model fidelity is a *two-way street*.



## Faithful Physical Model of Newton's Cradle?

- localized plastic deformation
- viscous damping
- acoustic wave propagation

If our goal is to understand macroscopic *system* behavior, then such a model may not be a good choice.

## Directly Modeling Discrete Phenomena

Difficulties that arise:

- signals are non-differentiable,
- chattering can occur around the discontinuity,
- Zeno behaviors can arise, and
- causality loops can arise.

Discrete physical phenomena are among the most poorly understood.

## Mixed Discrete and Continuous Models

Our approach uses:

- superdense model of time,
- generalized functions,
- modal models, and
- constructive semantics.

Detailed report is here:

<http://www.eecs.berkeley.edu/Pubs/TechRpts/2014/EECS-2014-15.html>

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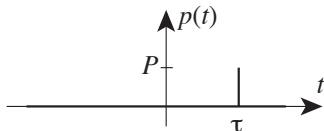
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## The Standard Model — Real-Valued Time

Let  $p: \mathbb{R} \rightarrow \mathbb{R}$  be a function that represents the momentum of the second ball in Newton's cradle, and let  $\tau$  be the time of the collision. Then

$$p(t) = \begin{cases} P & \text{if } t = \tau \\ 0 & \text{otherwise} \end{cases}$$

for some constant  $P$  and for all  $t \in \mathbb{R}$ .



## The Standard Model — Real-Valued Time

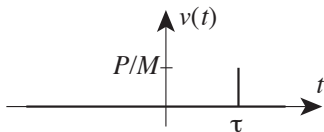
$$p(t) = \begin{cases} P & \text{if } t = \tau \\ 0 & \text{otherwise} \end{cases} \quad (1)$$

Momentum is proportional to velocity, so

$$p(t) = Mv(t),$$

where  $M$  is the mass of the ball. Hence, combining with (1),

$$v(t) = \begin{cases} P/M & \text{if } t = \tau \\ 0 & \text{otherwise.} \end{cases}$$



## The Standard Model — Real-Valued Time

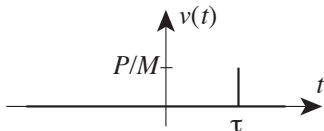
$$v(t) = \begin{cases} P/M & \text{if } t = \tau \\ 0 & \text{otherwise.} \end{cases} \quad (2)$$

The position of a mass is the integral of its velocity,

$$x(t) = x(0) + \int_0^t v(\tau) d\tau,$$

where  $x(0)$  is the initial position.

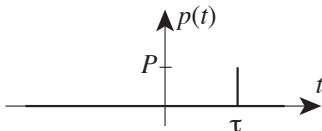
The integral is zero at all  $t$ ,  
so the ball does not move.



## The Standard Model — Real-Valued Time

$p: \mathbb{R} \rightarrow \mathbb{R}$ , where

$$p(t) = \begin{cases} P & \text{if } t = \tau \\ 0 & \text{otherwise} \end{cases}$$



**Problems with this model:**

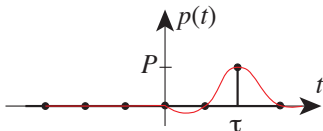
- Violates conservation of momentum at the time of the collision.
- The model cannot be unambiguously given a discrete representation.



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**Problems with this model:**

- Violates conservation of momentum at the time of the collision.
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## Superdense Time Model

Let  $p: \mathbb{R} \times \mathbb{N} \rightarrow \mathbb{R}$ , where

$$p(t, n) = \begin{cases} P & \text{if } t = \tau \wedge n = 1 \\ 0 & \text{otherwise} \end{cases}$$

I.e., at the time  $\tau$  of the collision,

- $p(\tau, 0) = 0$
- $p(\tau, 1) = P$
- $p(\tau, 2) = 0$



## Summary: Superdense Time

- **time stamp**:  $(t, n)$
- **model time**:  $t \in \mathbb{R}$
- **microstep** (or **index**):  $n \in \mathbb{N}$

Two time stamps  $(t_1, n_1)$  and  $(t_2, n_2)$  are

- **weakly simultaneous** if  $t_1 = t_2$
- **strongly simultaneous** if  $t_1 = t_2$  and  $n_1 = n_2$

Superdense time has been used by [Manna and Pnueli \(1993\)](#); [Maler et al. \(1992\)](#); [Lee and Zheng \(2005\)](#).

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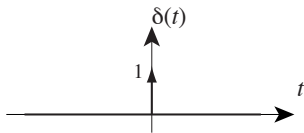
## Dirac Delta Function

The **Dirac delta function** is a function  $\delta: \mathbb{R} \rightarrow \mathbb{R}^+$  given by

$$\forall t \in \mathbb{R}, t \neq 0, \quad \delta(t) = 0, \quad \text{and}$$

$$\int_{-\infty}^{\infty} \delta(\tau) d\tau = 1.$$

That is, the signal value is zero everywhere except at  $t = 0$ , but its integral is unity.



## Impulsive Events with Dirac Delta Functions

Suppose  $x$  has a Dirac delta function with weight  $K$  occurring at  $t_1$  as follows,

$$x(t) = x_1(t) + K\delta(t - t_1),$$

where  $x_1$  is an ordinary continuous-time signal. Then

$$\int_{-\infty}^t x(\tau) d\tau = \begin{cases} \int_{-\infty}^t x_1(\tau) d\tau & t < t_1 \\ K + \int_{-\infty}^t x_1(\tau) d\tau & t \geq t_1 \end{cases}$$

The Dirac delta causes an instantaneous increment in the integral by  $K$  at time  $t = t_1$ .

## Direct Feedthrough

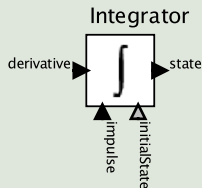
At time  $t$ , the *state* output is

$$x(t) = x_0 + \int_{t_0}^t \dot{x}(\tau) d\tau,$$

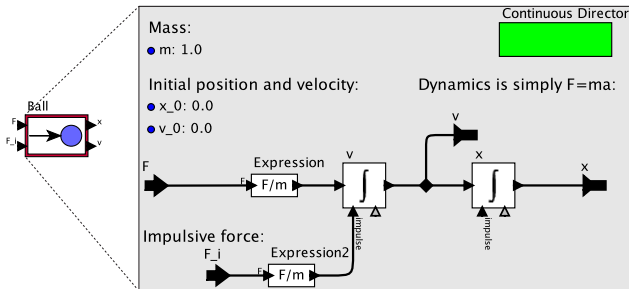
If the *impulse* input is present, then it adds immediately to  $x(t)$ .

The output at time  $t$  depends on the *impulse* input at time  $t$ , but not on the *derivative* input.

Ptolemy II Integrator has “impulse” input:



## Ball Model



The output  $v$  depends immediately on the input  $F_i$ , if it is present.



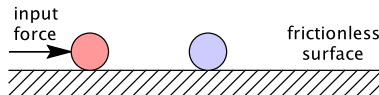
## Mixed Discrete and Continuous Models

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## Collisions



Conservation of momentum:

$$m_1 v'_1 + m_2 v'_2 = m_1 v_1 + m_2 v_2.$$

Conservation of kinetic energy:

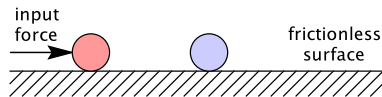
$$\frac{m_1 (v'_1)^2}{2} + \frac{m_2 (v'_2)^2}{2} = \frac{m_1 (v_1)^2}{2} + \frac{m_2 (v_2)^2}{2}.$$

We have two equations and two unknowns,  $v'_1$  and  $v'_2$ .

## After a Collision

Quadratic problem has two solutions.

**Solution 1:**  $v'_1 = v_1$ ,  $v'_2 = v_2$   
(ignore collision).

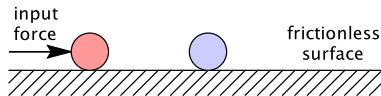


**Solution 2:**

$$v'_1 = \frac{v_1(m_1 - m_2) + 2m_2v_2}{m_1 + m_2}$$
$$v'_2 = \frac{v_2(m_2 - m_1) + 2m_1v_1}{m_1 + m_2}.$$

Note that if  $m_1 = m_2$ , then the two masses simply exchange velocities (Newton's cradle).

## Coefficient of Restitution

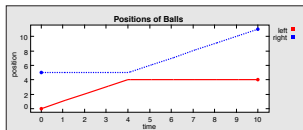


$$v_1' = \frac{em_2(v_2 - v_1) + m_1v_1 + m_2v_2}{m_1 + m_2} \quad (3)$$

$$v_2' = \frac{em_1(v_1 - v_2) + m_1v_1 + m_2v_2}{m_1 + m_2}. \quad (4)$$

The **coefficient of restitution**  $e$  is determined experimentally for a particular pair of materials and must lie in the range  $0 \leq e \leq 1$ .

## Collision Model

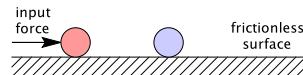
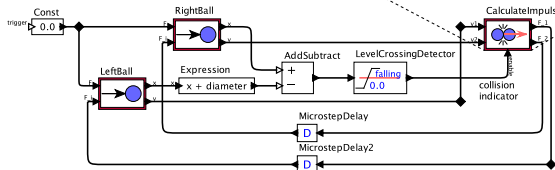
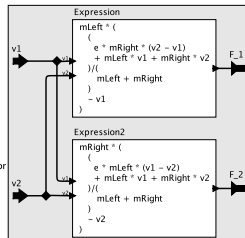


Model of two balls colliding.

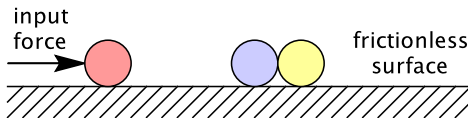
Parameters:

- mLeft: 1.0 left mass
- e: 1.0 coefficient of restitution
- mRight: 1.0 right mass
- diameter: 1.0

Continuous Director



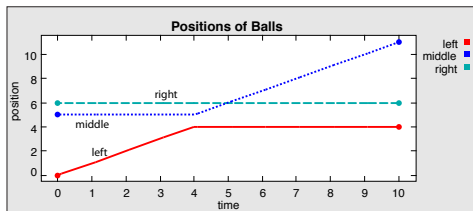
## Simultaneous, Causal Collisions



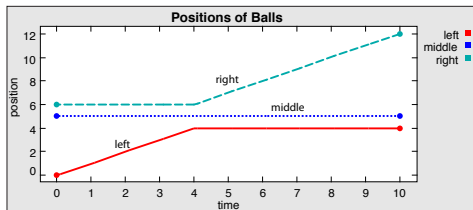
The middle ball instantly collides with the right ball, and the middle ball should not move.

## Collision Detection

Detecting discrete events is difficult. No zero crossing occurs. Failure to detect results in the upper plot.

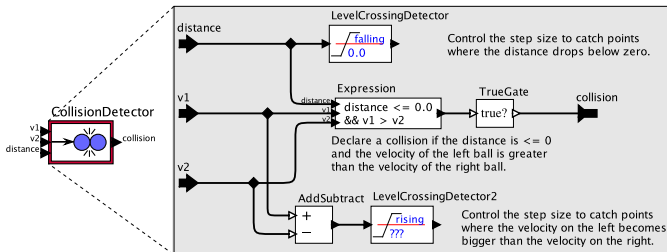


(a)



(b)

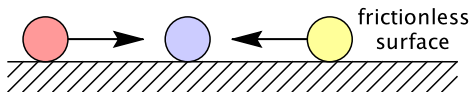
## Better Collision Detector



The LevelCrossingDetector actors regulate step sizes.



## Simultaneous, Noncausal Collisions

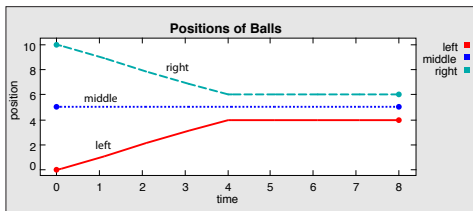


Simultaneous collisions where one collision does not cause the other.

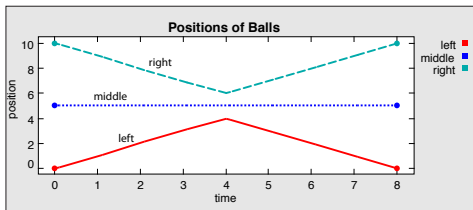
## Simultaneous Collisions

Newton's hypothesis  
superimposes impulsive  
forces, which cancel.

Poisson's hypothesis  
separates a compression  
phase and a restitution  
phase.



(a)

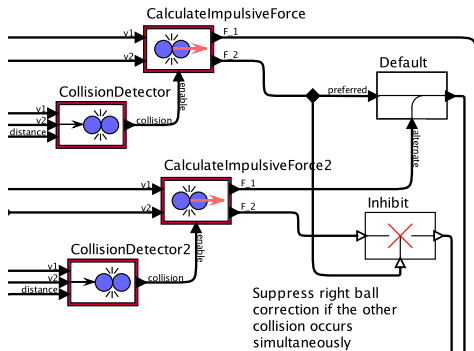


(b)

## Arbitrary Interleaving

Recall that one solution that conserves energy and momentum is to ignore a collision.

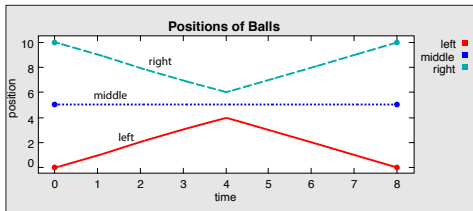
Here, if two collisions occur simultaneously, the right-middle collision is ignored.



CollisionSimultaneous.xml

## Arbitrary Interleaving

But then, one superdense index later, a second collision occurs, where the middle ball and the right ball collide, traveling towards each other.

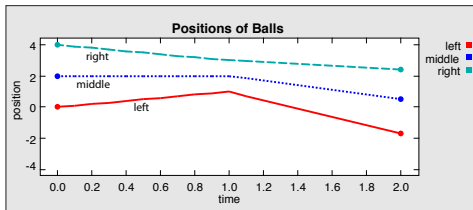


CollisionSimultaneous.xml

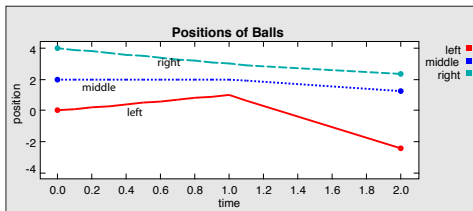
Result is like the Poisson solution.

## Arbitrary Interleaving Yields Nondeterminism

If the masses are different, the behavior depends on which collision is handled first!



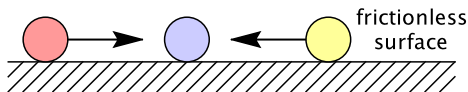
(a)



(b)

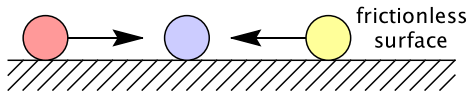
## Heisenberg Uncertainty Principle

We cannot simultaneously know the position and momentum of an object to arbitrary precision. But the reaction to these collisions depends on knowing position and momentum precisely.



## Heisenberg Uncertainty Principle

Arbitrary interleaving and nondeterministic results appears to be defensible on physical grounds.



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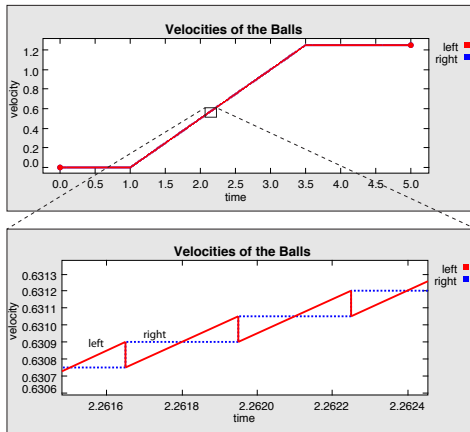
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## Pushing vs. Collisions

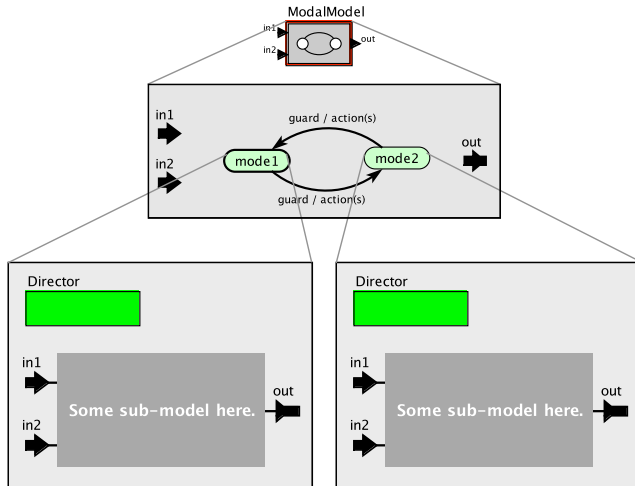
Consider two balls that start out touching, and from time 1 to 3.5 we apply a constant pushing force.

Treating these as collisions is very expensive.



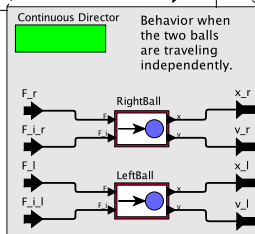
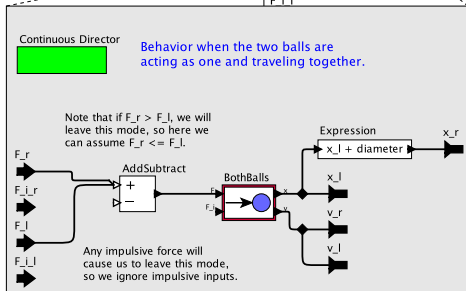
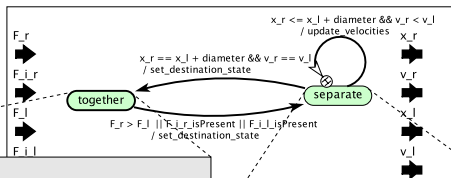
Collision5.xml

## Modal Models



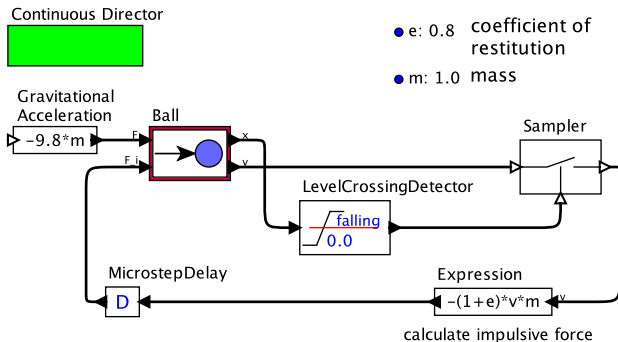
# Constructive Models of Discrete and Continuous Physical Phenomena

## Modal Pushing



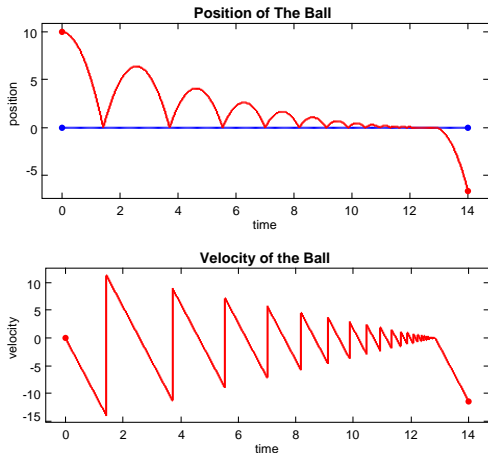
## Bouncing Ball — Zeno Conditions

Impulsive force  
is fed back to  
model the  
bounce.



## Zeno

Zero-crossing detection ultimately depends on numerical approximations to real numbers. It can only be done up to some precision.

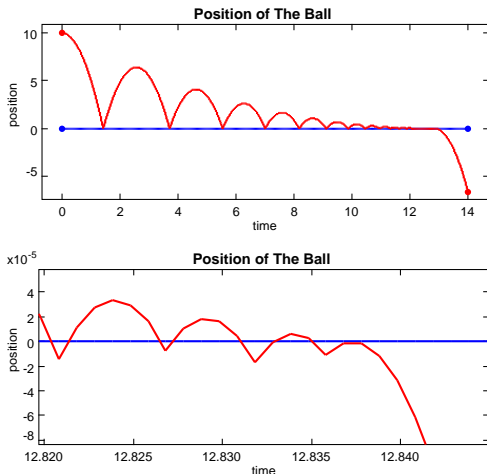


BouncingBall.xml

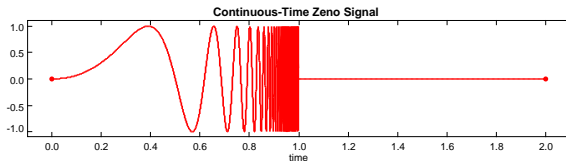
## Zeno

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Maybe tunneling is nature's way of resolving Zeno conditions.



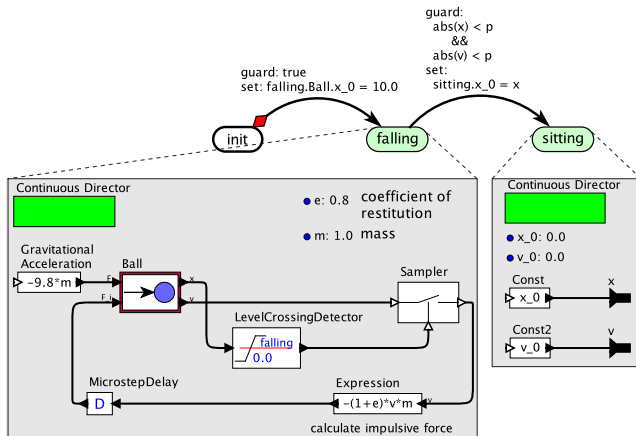
## Continuous Zeno Behavior



Zeno behaviors are not limited to discrete events.

## Modal Bouncing Ball

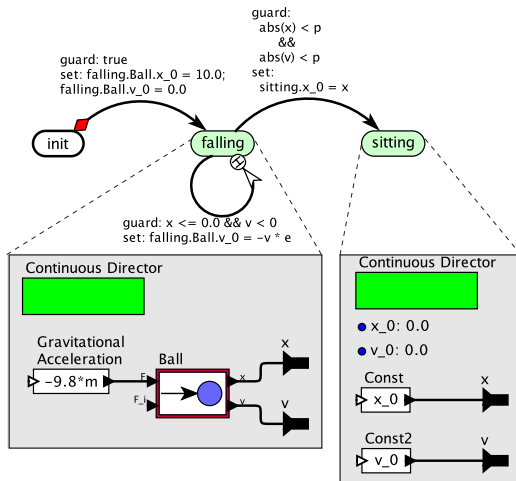
Impulsive force is fed back to model the bounce.





## Alternative Modal Bouncing Ball

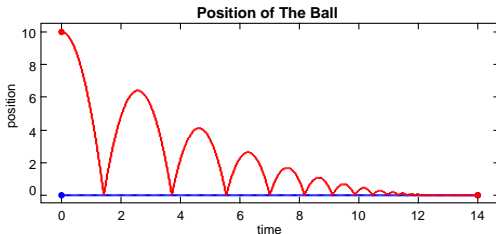
Mode transition handles the discrete change of energy and momentum.



BouncingBallModal2.xml

## Regimes of Models

Mode transition changes the model when the state of the system moves outside the regime over which the model is reasonable.



Note that *all* models have a limited regime over which they are reasonable.

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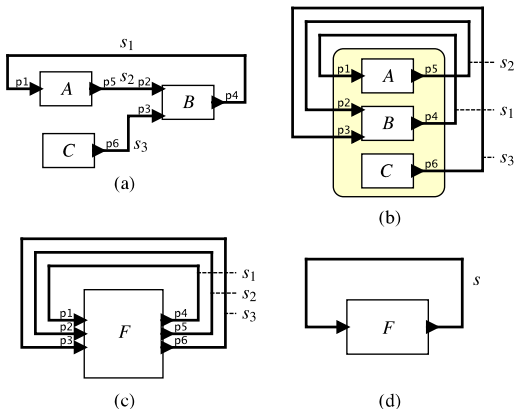
<http://www.eecs.berkeley.edu/Pubs/TechRpts/2014/EECS-2014-15.html>

## Constructive Fixed-Point Semantics

Every model reduces to a fixed point problem. For each  $t \in \mathbb{R} \times \mathbb{N}$ , find  $s(t)$  s.t.

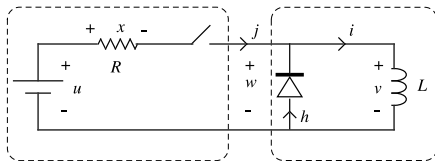
$$F(s(t)) = s(t).$$

With constructive semantics, solution is unique.  
See report.

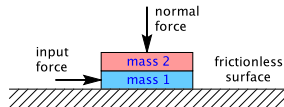


## Other Examples Considered in the Paper

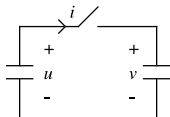
switches and diodes:



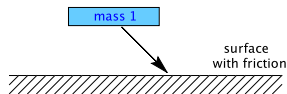
friction:



circuits with no solution:



collisions and friction:



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## References

- Lee, E. A., 2014: Constructive models of discrete and continuous physical phenomena. Tech. rep., EECS Department, UC Berkeley. Available from: <http://chess.eecs.berkeley.edu/pubs/1053.html>.
- Lee, E. A. and H. Zheng, 2005: Operational semantics of hybrid systems. In Morari, M. and L. Thiele, eds., *Hybrid Systems: Computation and Control (HSCC)*, Springer-Verlag, Zurich, Switzerland, vol. LNCS 3414, pp. 25–53. [doi:10.1007/978-3-540-31954-2\\_2](https://doi.org/10.1007/978-3-540-31954-2_2).
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**Questions?**